

Today :

- ① Fourier series
- ② some important PDEs.
- ③ Sol'n using Fourier series.

① Consider the interval  $[-L, L]$ .

consider complex valued (continuous, or Riemann integrable) function on  $[-L, L]$ .

Let  $f(x), g(x)$  be such nice functions on  $[-L, L]$ .  
we define "inner product"

$$\langle f, g \rangle = \frac{1}{2L} \int_{-L}^L \overline{f(x)} \cdot g(x) dx$$

[ compare this with Hermitian inner product on  $\mathbb{C}^N$ .  
 $\vec{x} = (x_1, \dots, x_N)$ ,  $\vec{y} = (y_1, \dots, y_N) \in \mathbb{C}^N$   
L  $\langle \vec{x}, \vec{y} \rangle = \sum_{i=1}^N \overline{x_i} \cdot y_i$

•  $e_n(x) := e^{\frac{\pi i x}{L} \cdot n}$  complex valued,  $n \in \mathbb{Z}$ .

[ it is  $2L$ -periodic,  $e_n(x+2L) = e_n(x)$ .

$$e_n(x+2L) = e^{\frac{\pi i}{L}(x+2L)n} = e^{\frac{\pi i}{L}x \cdot n + 2\pi i \cdot n}$$

L  $= e^{\frac{\pi i}{L} \cdot x \cdot n} \cdot e^{2\pi i n}$   
 $\swarrow$   
 $= 1$

$$\begin{aligned}
\langle e_n(x), e_m(x) \rangle &= \frac{1}{2L} \int_{-L}^L e^{\frac{\pi i}{L} x \cdot n} \cdot e^{\frac{\pi i}{L} x \cdot m} \cdot dx \\
&= \frac{1}{2L} \int_{-L}^L e^{\frac{\pi i}{L} x(m-n)} \cdot dx \\
&= \begin{cases} 1 & m=n. \\ 0 & m \neq n. \end{cases}
\end{aligned}$$

$k = m - n \neq 0$

$$\begin{aligned}
&\int_{-L}^L e^{\frac{\pi i}{L} \cdot k \cdot x} dx = \left. \frac{e^{\frac{\pi i}{L} \cdot k \cdot x}}{\frac{\pi i}{L} \cdot k} \right|_{-L}^L \\
&= \frac{e^{\frac{\pi i}{L} \cdot L \cdot k} - e^{\frac{\pi i}{L} (-L) k}}{\frac{\pi i}{L} \cdot k} = \frac{e^{k\pi i} - e^{-k\pi i}}{\frac{\pi i}{L} \cdot k} \\
&= \frac{(-1)^k - (-1)^{-k}}{\frac{\pi i}{L} \cdot k} = 0. \quad e^{\pi i} = (-1)
\end{aligned}$$

remark: if we consider the " $L^2$  space on  $[-L, L]$ " with the above inner product, then  $\{e_n(x)\}_{n \in \mathbb{Z}}$  forms an orthonormal basis.

Given any continuous function  $f(x)$  on  $[-L, L]$ , we can define its complex Fourier coefficients  $\{c_n\}_{n \in \mathbb{Z}}$ .

$$c_n := \langle e_n, f \rangle = \frac{1}{2L} \int_{-L}^L e^{-\frac{\pi i}{L} \cdot x \cdot n} \cdot f(x) dx.$$

Complex Fourier series :

$$f(x) \stackrel{?}{=} \sum_{n \in \mathbb{Z}} C_n \cdot e_n(x).$$

Compare : suppose  $V$  is a finite dim, Hermitian vector space  
and  $\{e_1, \dots, e_n\}$  is an ONB.

Then, for any  $v \in V$ , we have  
$$v = \sum_{i=1}^n c_i \cdot e_i, \quad \text{and the coeff}$$

$$c_i = \langle e_i, v \rangle.$$

- For  $f(x)$  continuous, then for any  $x \in [-L, L]$ , we have

$$f(x) = \sum_{n \in \mathbb{Z}} e_n(x) \cdot \langle e_n, f \rangle = \lim_{N \rightarrow \infty} \sum_{n=-N}^N e_n(x) \langle e_n, f \rangle.$$

- Real functions's Fourier expansion.

Given a real valued function  $f(x)$  on  $[-L, L]$ ,  
we can use

$$1, \sin\left(\frac{\pi x}{L} \cdot n\right), \cos\left(\frac{\pi x}{L} \cdot n\right), \quad n=1, 2, \dots$$

to expand  $f$ . i.e.

$$f(x) = a_0 + \sum_{n=1}^{\infty} \left( a_n \cdot \cos\left(\frac{\pi x}{L} \cdot n\right) + b_n \cdot \sin\left(\frac{\pi x}{L} \cdot n\right) \right).$$

These  $\overset{S_n(x)}{\parallel} \sin\left(\frac{\pi x}{L} \cdot n\right)$  and  $\overset{C_n(x)}{\parallel} \cos\left(\frac{\pi x}{L} \cdot n\right)$   $\swarrow \frac{1}{\text{forms}}$

an ONB.

$$\begin{aligned} e_n(x) &= e^{\frac{\pi i}{L} x \cdot n} = \cos\left(\frac{\pi}{L} \cdot x \cdot n\right) + i \sin\left(\frac{\pi x}{L} \cdot n\right) \\ &= C_n(x) + i \cdot S_n(x). \end{aligned}$$

$$\begin{aligned} e_{-n}(x) &= C_{-n}(x) + i S_{-n}(x) \\ &= C_n(x) - i S_n(x). \end{aligned}$$

$$C_n(x) = \frac{e_n + e_{-n}}{2}$$

$$S_n(x) = \frac{e_n - e_{-n}}{2i}$$

$$\begin{aligned} \langle C_n, S_n \rangle &= \left\langle \frac{e_n + e_{-n}}{2}, \frac{e_n - e_{-n}}{2i} \right\rangle \\ &= \frac{1}{4i} \left( \langle e_n, e_n \rangle - \langle e_{-n}, e_{-n} \rangle \right) \\ &= \frac{1}{4i} (1 - 1) = 0. \end{aligned}$$

$$\begin{aligned} \langle C_n, C_n \rangle &= \left\langle \frac{e_n + e_{-n}}{2}, \frac{e_n + e_{-n}}{2} \right\rangle \\ &= \frac{1}{4} \cdot (\langle e_n, e_n \rangle + \langle e_{-n}, e_{-n} \rangle) \\ &= \frac{1}{4} \cdot 2 = \frac{1}{2}. \end{aligned}$$



② Famous <sup>partial</sup> Diff Eq.

Given a function that depends on more than one variable. say  $f(x, y)$ ,

we can talk about its partial derivative.

$$\partial_x f(x_0, y_0) := \lim_{\varepsilon \rightarrow 0} \frac{f(x_0 + \varepsilon, y_0) - f(x_0, y_0)}{\varepsilon}.$$

$$\partial_y f(x_0, y_0) := \lim_{\varepsilon \rightarrow 0} \frac{f(x_0, y_0 + \varepsilon) - f(x_0, y_0)}{\varepsilon}.$$

we can define higher order partial derivatives

$$\partial_x(\partial_x f), \quad \partial_y(\partial_x f),$$

$$\partial_x(\partial_y f), \quad \partial_y(\partial_y f).$$

In the case that  $f$  is "twice differentiable"

then.  $\partial_x \partial_y f = \partial_y \partial_x f.$

Ex:  $f(x, y) = x \cdot y.$  then

$$\partial_x f = y, \quad \partial_y f = x.$$

$$\partial_x \partial_x f = \partial_x (y) = 0, \quad \partial_y \partial_y f = 0.$$

$$\partial_x \partial_y f = \partial_x (x) = 1.$$

$$\partial_y \partial_x f = \partial_y (y) = 1.$$

• We will use notation.  $f_x = \partial_x f$ ,  $f_{xx} = \partial_x \partial_x f = \partial_x^2 f$ .

• Laplacian :  $\Delta$ .

on  $\mathbb{R}^n$   $\Delta = (\partial_{x_1})^2 + (\partial_{x_2})^2 + \dots + (\partial_{x_n})^2$ .

•  $\Delta f$  measures how much a function "bend".

Ex: on  $\mathbb{R}$ .  $\Delta f(x) = \partial_x^2 f(x)$ .

$$\Delta f(x) = 0 \iff \partial_x^2 f(x) = 0$$

$$\iff f(x) = a + bx. \quad \text{straight line.}$$

Ex: on  $\mathbb{R}^2$ ,  $\Delta f(x,y) = \partial_x^2 f + \partial_y^2 f$ .

$$\Delta f(x,y) = 0 \iff \partial_x^2 f(x,y) + \partial_y^2 f(x,y) = 0$$

ex sol'n:  $f(x,y) = a + bx + cy$ .

$$f(x,y) = xy$$

$$\text{or } f(x,y) = x^2 - y^2,$$

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 The sol'n to  $\Delta f(x,y) = 0$  <sup>are</sup> ~~is~~ called  
 harmonic functions.

• Eigenvalue problem for  $\Delta$ :

$$\begin{array}{l} \text{Ex: let } \Omega \subset \mathbb{R}^n, \quad f: \Omega \rightarrow \mathbb{R}. \\ \left\{ \begin{array}{l} \Delta f = \lambda f \\ f|_{\partial\Omega} = 0. \end{array} \right. \quad \left| \begin{array}{l} \lambda : \text{eigenvalue} \\ \text{for } \Delta. \\ f : \text{eigenvector.} \end{array} \right. \end{array}$$

$$\text{Ex: } [0,1] \subset \mathbb{R}. \quad f(x) \quad \text{s.t.}$$

$$\left\{ \begin{array}{l} \partial_x^2 f = \lambda f \\ f(0) = 0, \quad f(1) = 0 \end{array} \right.$$



in 1-dim, we have the sequence of sol'n.

$$\left\{ \begin{array}{l} f_n(x) = \sin(\pi n \cdot x) \\ \lambda_n = -(\pi n)^2. \end{array} \right.$$

$$n = 1, 2, \dots$$

$$\partial_x f_n(x) = \partial_x \sin(\pi n \cdot x) = \cos(\pi n \cdot x) \cdot (\pi n)$$

$$\partial_x^2 f_n(x) = \partial_x \cos(\pi n \cdot x) \cdot (\pi n) = -\sin(\pi n \cdot x) \cdot (\pi n)^2$$

• Heat equation:

$$\partial_t f(t,x) = \Delta_x f(t,x)$$

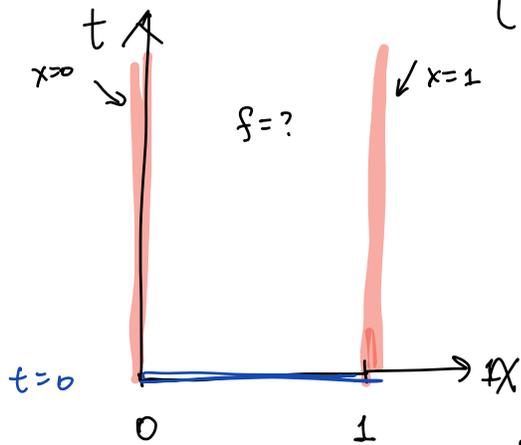
$$f: \begin{array}{l} \mathbb{R} \times \mathbb{R}^n \rightarrow \mathbb{R}. \\ t \quad x \end{array}$$

in special case:

want  $f: [0, \infty)_t \times [0, 1]_x \rightarrow \mathbb{R}$ .

such that.

$$\begin{cases} \partial_t f(t, x) = \partial_x^2 f(t, x) \\ f(0, x) = f_0(x) \\ f(t, 0) = 0, \quad f(t, 1) = 0 \end{cases}$$



Fourier expansion.

If  $f_0(x) = \sum_{n=1}^{\infty} b_n \cdot \sin(\pi x \cdot n)$ .

then we know how each  $\sin(\pi x \cdot n)$  evolve.  
 i.e. if  $f_0(x) = \sin(\pi x \cdot n)$ ,  
 then  $f(t, x) = e^{-(\pi n)^2 \cdot t} \cdot \sin(\pi x \cdot n)$ .  
 $\partial_t f = -(\pi n)^2 \cdot e^{-(\pi n)^2 \cdot t} \cdot \sin(\pi x \cdot n)$   
 $\partial_{xx} f = e^{-(\pi n)^2 \cdot t} \cdot (-\pi n)^2 \cdot \sin(\pi x \cdot n)$   
 $\Rightarrow \partial_t f = \partial_{xx} f$ .

then  $f(t, x) = \sum_{n=1}^{\infty} b_n \cdot e^{-(\pi n)^2 \cdot t} \cdot \sin(\pi x \cdot n)$ .